# Finite size phase transitions in QCD with adjoint fermions

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## Summary of the talk



### Motivation and background

- Introduction Large N<sub>C</sub> theories
- QCD with Adjoint fermions on the lattice



- Phase space diagram Chiral limit
- Chiral properties
- Summary and outlook

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Introduction - Large N<sub>C</sub> theories QCD with Adjoint fermions on the lattice

## Outline



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## Introduction

Paper: G.Cossu, M.D'Elia arXiv:0904.1353 [hep-lat]

#### Large $N_c$ theories

- Useful tool for the study of non perturbative QCD
- QCD with fermions in the antisymmetric (AS) representation is a different  $N_c \rightarrow \infty$  limit of ordinary QCD (equal at  $N_c = 3$ ).
- Orientifold equivalence ⇒ QCD(AS) = QCD(Adj) in this limit (Shifman, Armoni, Veneziano, Patella)
- Kovtun, Unsal, Yaffe claims that QCD(Adj) stays confined in a ℝ<sup>3</sup> × S<sub>1</sub> space if we shrink the compact dimension. Easier calculations extrapolated back to ordinary QCD by the various equivalences.

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## QCD(Adj) at small volume - 1

Kovtun, Unsal, Yaffe JHEP, 0706 (2007) 019 Analysis of 1-loop effective potential for the Polyakov Line along the compact dimension ( $L_c$ ),  $N_f$  Majorana fermions. Gives all the relevant information at small  $L_c$ .

#### Polyakov Loop Potential

$$V_{\rm eff}(\Omega) = \frac{1 \mp N_f}{24\pi^2 L_c^4} \sum_{i,j=1}^{N_c} [\phi_i - \phi_j]^2 ([\phi_i - \phi_j] - 2\pi)^2$$
(1)

+/- for periodic/antiperiodic boundary conditions (PBC/ABC)  $e^{\phi_i}$  eigenvalues of the Polyakov Loop.

- ABC or  $N_f = 0$  minimized if eigvals coincide
- PBC: repulsive potential for eigvals, disordered.

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## QCD(Adj) at small volume - 2

- Results confimed by 1-loop lattice calculations
- Same conclusions in a  $SU(N_C)$  + trace deformation: i.e. small and large volumes are analytically connected.
- The essential requirement is that no phase transition occurs: we will address this point by numerical lattice simulations of QCD with 2 Dirac adjoint fermions and  $N_C = 3$ .
- We will discuss the chiral and continuum limit.
- We conclude that the claim by Kovtun, Unsal and Yaffe is highly non trivial and should be considered with care.

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Introduction - Large  $N_C$  theories QCD with Adjoint fermions on the lattice

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## Lattice discretization of QCD(Adj)

#### **Adjoint Fermions**

- 8 color d.o.f.  $Q = Q^a \lambda_a (\lambda_a \text{ Gell-Mann's matrices})$
- $U^{ab}_{(8)\ i,\mu} = \frac{1}{2} \operatorname{Tr} \left( \lambda^a U_{(3)\ i,\mu} \lambda^b U^{\dagger}_{(3)\ i,\mu} \right)$  (8-dim rep. of links)
- $S = S_G [U_{(3)}] + \sum_{i,j,a,b} \bar{Q}_i^a M [U_{(8)}]_{i,j}^{a,b} Q_j^b$  (Action)

Staggered fermions

#### Partition Function

• 
$$Z = \int \mathcal{D}U\mathcal{D}\Phi_e e^{-S_G} \exp\left(-\Phi_e^t (M^t M)_{ee}^{-1} \Phi_e\right)$$

 Integral sampled by standard Hybrid Monte Carlo Φ algorithm.

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## Center symmetry and Polyakov Loop

#### Center symmetry

YM theories are invariant if we multiply all links in a temporal slice by a center element. Fermionic terms (in the fundamental representation) break this symmetry

#### Polyakov Loop

$$L_3 = \operatorname{Tr} \prod_{t=0}^{t=N_t-1} U_3(\vec{x},t)$$

(2)

- The action of QCD(Adj) is invariant for center symmetry.
- The Fundamental Polyakov Loop is not and therefore is an order parameter for its spontaneous breaking.

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## Simulations parameters - 1

- ABC : thermodynamical system.
- PBC : no thermodynamical interpretation although we will use the same terminology.
- We simulated the variation of length of the compactified dimension by varying β (gauge coupling) and keeping the number of sites constant.
- By asymptotic freedom lattice spacing  $a \rightarrow 0$  as  $\beta \rightarrow \infty$ .
- At infinite mass (pure YM) we have two distinct phases separated by a first order transition.
- At the chiral point we have two confined phases at the opposite limits of compact dimension (*L<sub>c</sub>* → 0, ∞).
- A highly non trivial phase structure must be found at intermediate masses.

## Simulations parameters - 2

- Several simulations with different masses toward the chiral limit are essential (*am* = 0.50, 0.10, 0.05, 0.02, 0.01).
- Lattice dimensions  $16^3 \times L_c$  with  $L_c = 4$  for most of the simulations.

#### **Technicalities**

- Trajectory length: 0.5
- Integration step δt from 0.025 to 0.005 depending on the mass to keep acceptance rate around 80%
- MC data from 2k to 20k depending on paramters
- Performed on the APEmille installation in Pisa and the apeNEXT facility in Rome.

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Phase space diagram - Chiral limit Chiral properties Summary and outlook

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## Phase transitions am= 0.02



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## Phase transitions - Cumulative



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## Phase transitions - Continuum Limit $16^3 \times 6$



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## Phase transitions - Fixed lattice spacing



G.Cossu Finite size

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## Chiral limit extrapolations



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Phase space diagram - Chiral limit

#### Chiral properties

Summary and outlook

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## Chiral condensate extrapolations



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## Relation with center symmetry - 1

• Bilgici, Bruckmann, Gattringer, Hagen proposed a new operator to detect confinement that is related to the chiral properties.

$$\Sigma_{1} = -\frac{1}{2\pi} \int_{0}^{2\pi} d\varphi \; e^{-i\varphi} \langle \bar{\psi}\psi \rangle_{m}^{(\varphi)} \tag{3}$$

where  $\langle \bar{\psi}\psi \rangle_m$  is the chiral condensate measured with an assigned phase for the fermionic boundary conditions in the compactified dimension. In pure YM this reduces to ordinary Polyakov Loop.

• We measured an approximation :  $(\langle \bar{\psi}\psi \rangle_m^{(0)} - \langle \bar{\psi}\psi \rangle_m^{(\pi)})$ 

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## Relation with center symmetry - 2



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## Conclusions

- We have shown that the claim by Kovtun, Unsal and Yaffe must be taken with care.
- At least two phase transitions seem to occur at small masses for PBC QCD(Adj) if we shrink the volume toward small compactification radii.
- A new phase appear at intermediate masses (it is not clear if it survives the chiral limit)
- We verified that center symmetry is reflected in the dependence of the chiral condensate on BC.
- Chiral limit must be further investigated and different discretization methods would be useful to detect systematic errors.

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