# Spectral Curves and Localization in Random Non-Hermitian Tridiagonal Matrices 

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## Plan of the talk

(9) Introduction
(2) Spectral properties
(3) Hatano-Nelson deformation

4 Spectral duality and other wonders
(5) Conclusions

## Introduction

## Random tridiagonal matrices

The hermitian case is very well studied; a lot is known on:

- spectral properties
- localization of eigenvectors
one major application being the Anderson model $(d=1)$


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- spectral properties
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one major application being the Anderson model ( $d=1$ )

We made a study of non-Hermitian tridiagonal random matrices with corners, which find applications in

- systems with asymmetric hopping amplitudes
- 1D random walks
- models for biological systems
we consider matrices of the form

$$
M=\left[\begin{array}{cccc}
a_{1} & b_{1} & & c_{1} \\
c_{2} & \ddots & \ddots & \\
& \ddots & \ddots & b_{n-1} \\
b_{n} & & c_{n} & a_{n}
\end{array}\right]
$$

where the elements $a, b, c$ are random complex numbers. We take them to be uniformly distributed in the unitary disk. We are interested in studying the large $n$ limit.
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spectrum of a sample matrix for $n=800$

Let us look at the eigenvalue equation, which takes the form (corners account for b.c.)

$$
c_{k} u_{k-1}+a_{k} u_{k}+b_{k} u_{k+1}=E u_{k}
$$

where $E$ is complex; it is quite easy to show that

however, we empirically find $|E| \leq 2$

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Eigenvalue density for matrices of the $M$ ensemble, $n=1000,100$ samples

The form of the eigenvalue equation makes it natural to introduce a $2 \times 2$ transfer matrix

$$
t(E)=\prod_{k=1}^{n}\left[\begin{array}{cc}
b_{k}^{-1}\left(E-a_{k}\right) & -b_{k}^{-1} c_{k} \\
1 & 0
\end{array}\right]
$$

such that, if

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\Psi_{n}=\binom{u_{n+1}}{u_{n}}
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the two eigenvalues of $t(E)$ can be written as

$$
z_{ \pm}^{n}=e^{n\left(\xi_{ \pm}+i \varphi_{ \pm}\right)}
$$

with the interesting property that, for large $n$ :

$$
\xi_{+}+\xi_{-}=\frac{1}{n} \log |\operatorname{det} t(E)|=\frac{1}{n} \sum_{k=1}^{n}\left(\log \left|c_{k}\right|-\log \left|b_{k}\right|\right) \rightarrow 0
$$

Apply to $t(E)$ the theory of random matrix products, for large $n$ :

- the positive exponent $\xi_{+}$is independent of $n$ and the realization of randomness
- and converges to the Lyapun ov exponent of the matrix ensemble
- given by an extension of the Thouless formula to non-Hermitian matrices:

$$
\gamma(E)=\int d^{2} E^{\prime} \rho_{0}\left(E^{\prime}\right) \log \left|E-E^{\prime}\right|-\langle\log | b| \rangle
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where $\rho_{0}$ is the eigenvalue density of the ensemble of matrices $M(\langle\log | b\rangle=-1 / 2$ in our case $)$

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## The Lyapunov exponent is related to the localization properties of eigenvectors, since asimptotically

$$
\left|u_{k}(E)\right| \simeq \exp \left(-\gamma(E)\left|k-k_{0}(E)\right|\right)
$$

giving to $\gamma(E)$ the meaning of an inverse localization length associated to the eigenstate with eigenvalue $E$


Numerical evaluation of the Lyapunov exponent for the $M$ ensemble, $n=1000,100$ samples

Let us look again at

$$
\Psi_{n}=t(E) \Psi_{1}
$$

What if we insist that $\Psi_{1}$ is an eigenvector of $t(E)$ ?

$$
t(E) \Psi_{1} \equiv z^{n} \Psi_{1}=\Psi_{n}
$$

this is clearly equivalent to a special deformation of the boundary conditions

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u_{n+1}=z^{n} u_{1}, u_{0}=z^{-n} u_{n}
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Following Hatano and Nelson, we introduce the following deformation of our matrices by a complex parameter $z=e^{\xi+i \varphi}$ :

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M\left(z^{n}\right)=\left[\begin{array}{cccc}
a_{1} & b_{1} & & z^{n} c_{1} \\
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the deformed matrix is related to

$$
M_{b}(z)=\left[\begin{array}{cccc}
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by a similarity transformation, $M_{b}(z)=S M\left(z^{n}\right) S^{-1}$, through a diagonal matrix with entries $S_{i i}=z^{i}$.
The two deformed matrices have the same spectrum, but $M_{b}$ is more convenient for numerical investigations.
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Eigenvalue spectrum with a HN deformation ( $\xi=0.5, \varphi=0$ ), compared with the undeformed case ( $n=800$ ).
A "bubble" has appeared... The eigenvalues outside it have not moved



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By changing $\xi$ to 0.7 (left) and 1.0 (right) $(\varphi=0)$ the bubble grows and sweeps the entire spectrum

change $\varphi$ with a step $2 \pi / n \ldots$ with $\xi=0.5$ (left) and $\xi=1$ (right)... the eigenvalues out of the bubble do not move... while those on the bubble shift and stay on it

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## We now wish to explain all these observed features.

- let $u(E)$ be an eigenstate of $M\left(z^{n}\right)$
- this implies that $S u$ is an eigenstate of $M_{b}(z)$
- given the form of $S$ and the localization properties of the eigenvectors $u$, one expects for large $n$ :
$\left|(S u)_{k}\right| \simeq \exp ((\xi-\gamma(E)) k)$
implying localization for $\gamma(E)>\xi$ and delocalization for
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Proposition: in the large $n$ limit, $M_{b}\left(e^{\xi+i \varphi}\right)$ has no eigenvalues inside the circle of radius $R$, where

$$
\gamma(R)=\xi
$$

we will prove this statement with

## - the Argument Principle

- a spectral duality

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- the Argument Principle
- a spectral duality
by the Argument Principle the number of zeroes of the analytic function

$$
f(E)=\operatorname{det}\left(E-M_{b}(z)\right)
$$

inside a disk is given by the variation of

$$
\frac{\arg (f(E))}{2 \pi}
$$

along the contour of the disk
on the other hand, the following duality identity holds

$$
\operatorname{det}\left[E-M_{b}(z)\right]=-\frac{1}{z^{n}}\left(b_{1} \cdots b_{n}\right) \operatorname{det}\left[t(E)-z^{n}\right]
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## we therefore have


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we therefore have

$$
\begin{aligned}
& \arg \operatorname{det}\left[E-M_{b}\left(e^{\xi+i \varphi}\right)\right]=\text { const. } \\
& \quad+\arg \left[e^{n\left(\xi_{+}-\xi\right)+i n\left(\varphi_{+}-\varphi\right)}-1\right] \\
& \quad+\arg \left[e^{n\left(\xi_{-}-\xi\right)+i n\left(\varphi_{-}-\varphi\right)}-1\right]
\end{aligned}
$$

# let us now put things together 

- $\xi>0$
- $\xi_{+}>0$
- in the large $n$ limit

We also identify $\xi_{+}(E)$ with the Lyapunov exponent $\gamma(|E|)$. The variation of $\arg f(E) / 2 \pi$ along a circumference of radius $R$ is zero if $\xi>\gamma(R)$. $\square$
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- we study the spectral properties of complex non-hermitian tridiagonal matrices with corners in the large $n$ limit
- we employ a Hatano-Nelson deformation to open a hole in the spectrum of our matrices
- the eigenvalues swept to the boundary of the hole correspond to states that are no longer Anderson localized
- this is explained in terms of a spectral duality, stability of the Lyapunov exponent and the Argument Principle
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